Introduction to Special Relativity



Manuel Franco Sevilla
University of Maryland

11th August 2025

Quarknet @ UMD

summer workshop







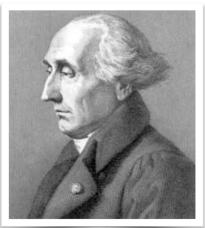


Physics complete in the XIX century?

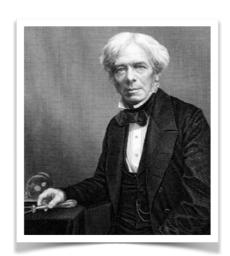


- ~ From the XVII to the XIX centuries, extraordinary progress in our fundamental understanding of the universe
 - → Newton's Gravitation, Classical and Statistical Mechanics, Thermodynamics, Electromagnetism

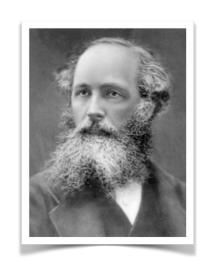












$$F = G \frac{Mm}{R^2}$$

$$oldsymbol{
abla} \cdot oldsymbol{E} = rac{
ho}{\epsilon_0}$$

$$\nabla \cdot \boldsymbol{B} = 0$$

$$oldsymbol{
abla} imes oldsymbol{E} = -rac{\partial oldsymbol{E}}{\partial t}$$

$$\nabla \cdot \boldsymbol{E} = \frac{\rho}{\epsilon_0}$$
 $\nabla \cdot \boldsymbol{B} = 0$ $\nabla \times \boldsymbol{E} = -\frac{\partial \boldsymbol{B}}{\partial t}$ $\nabla \times \boldsymbol{B} = \mu_0 \left(\boldsymbol{J} + \frac{\partial \boldsymbol{E}}{\partial t} \right)$

Feeling at the end of the XIX century

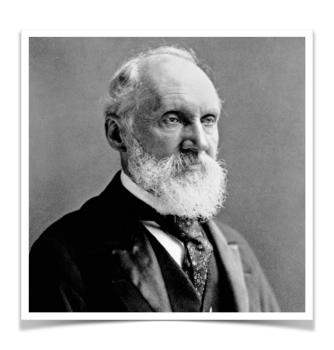
"There is nothing new to be discovered in physics now, All that remains is more and more precise measurements."



Lord Kelvin's clouds



"The beauty and clearness of the dynamical theory, which asserts heat and light to be modes of motion, is at present obscured by two clouds"



THE

LONDON, EDINBURGH, AND DUBLIN

PHILOSOPHICAL MAGAZINE

AND

JOURNAL OF SCIENCE

[SIXTH SERIES.]

JULY 1901.

I. Nineteenth Century Clouds over the Dynamical Theory of Heat and Light *. By The Right. Hon. Lord Kelvin, G.C.V.O., D.C.L., LL.D., F.R.S., M.R.I. †.







Michelson-Morley experiment



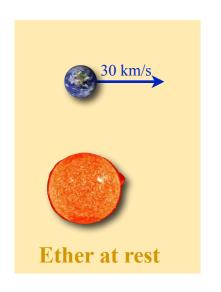
Maxwell equations predicted electromagnetic waves always traveling at speed $c=1/\sqrt{\mu_0 \epsilon_0}$

Speed c with respect to what???

$$\nabla^2 \mathbf{E} - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \mathbf{E} = 0$$

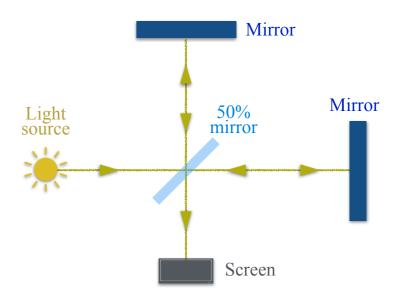
Hypothesis was that Earth was moving at least at 30 km/s with respect to Ether

If light moved at c with respect to Ether, beam would take different times when aligned with Earth's speed versus transverse to it





Michelson invented in 1881 (Naval academy, Annapolis) an **interferometer** that could measure a difference of one part in $c/v \approx 10^4$



No evidence for light traveling at different speeds in different directions

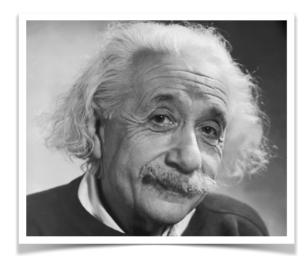
Albert Michelson would go onto win the Nobel prize for Physics in 1907, the first American to do so!





Einstein's postulates





Relativity principle: all laws of Physics have the same form in all inertial frames of reference

The speed of light in vacuum has the same value c in every direction in all inertial frames of reference

Zur Elektrodynamik bewegter Körper; von A. Einstein.

dies für die Größen erster Ordnung bereits erwiesen ist. Wir wollen diese Vermutung (deren Inhalt im folgenden "Prinzip der Relativität" genannt werden wird) zur Voraussetzung erheben und außerdem die mit ihm nur scheinbar unverträgliche Voraussetzung einführen, daß sich das Licht im leeren Raume stets mit einer bestimmten, vom Bewegungszustande des emittierenden Körpers unabhängigen Geschwindigkeit V fortpflanze. Diese beiden Voraussetzungen genügen, um zu einer einfachen und widerspruchsfreien Elektrodynamik bewegter Körper zu gelangen unter Zugrundelegung der Maxwellschen Theorie für ruhende Körper. Die Einführung eines "Lichtäthers" wird sich

On the Electrodynamics of Moving Bodies; A. Einstein

equations of mechanics hold good. We will raise this conjecture (the purport of which will hereafter be called the "Principle of Relativity") to the status of a postulate, and also introduce another postulate, which is only apparently irreconcilable with the former, namely, that light is always propagated in empty space with a definite velocity c which is independent of the state of motion of the emitting body. These two postulates suffice for the attainment of a simple and consistent theory of the electrodynamics of moving bodies based on Maxwell's theory for stationary bodies. The introduction of a "luminiferous ether" will

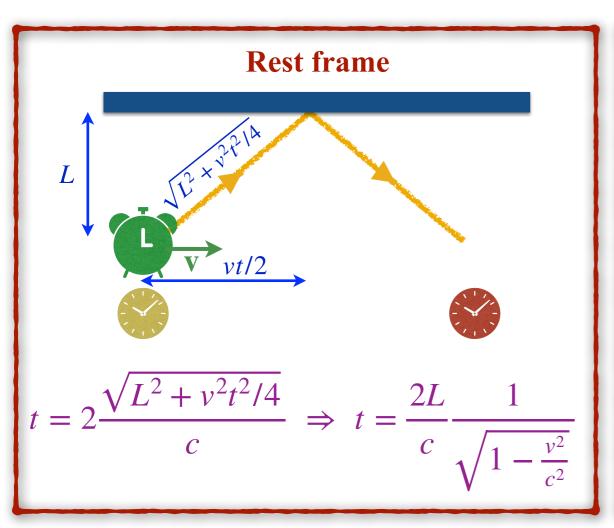
Einstein's 1905 paper is considered the first complete account of Special relativity, but many elements had also been derived by Hendrik Lorentz and Henri Poincaré

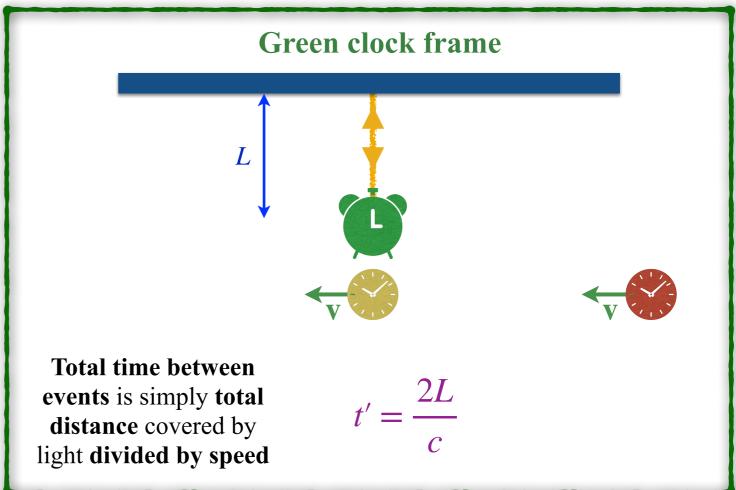


Time dilation



~ What's the time passed between light is emitted and absorbed by the green clock?





We define a new quantity γ

$$\gamma \equiv \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$t=t'\gamma$$

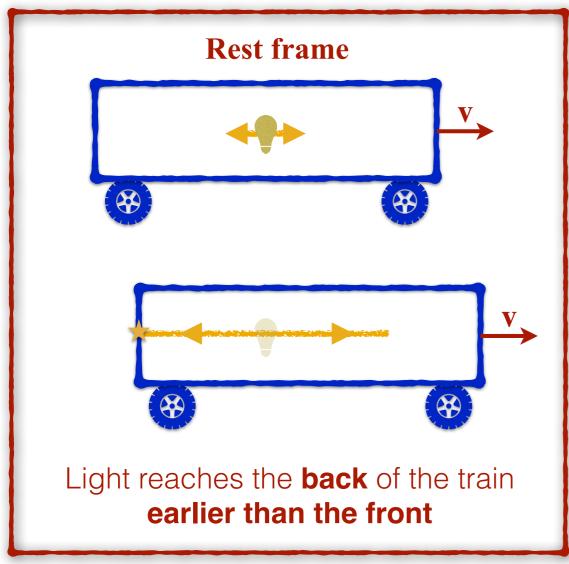
Based on Einstein postulates alone, time passes more slowly when moving

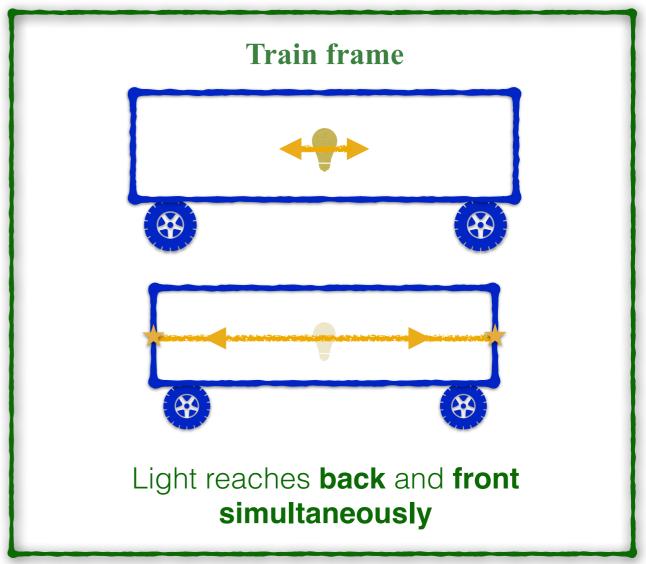


Simultaneity is frame dependent



- ~ A light goes off in the middle of a moving train
 - → Do the light rays reach the front/back of the train at the same time?



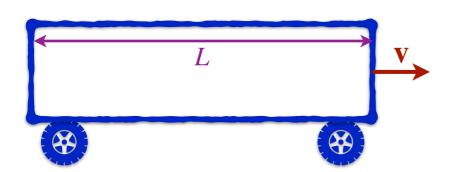


Whether two events that occurred far apart are simultaneous or not depends on the frame of reference



Length affected too?



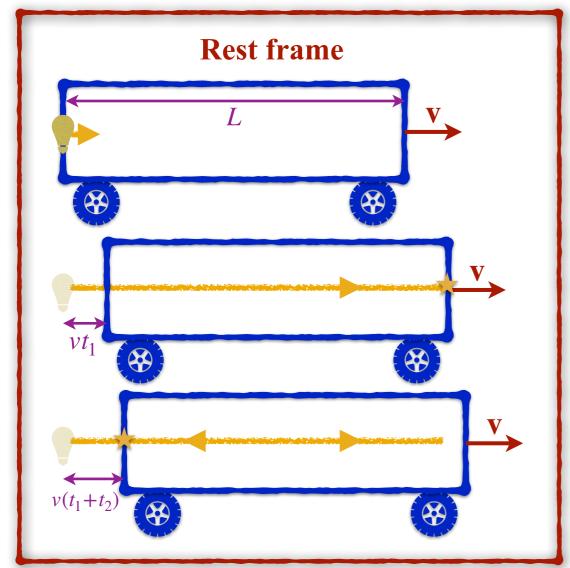


→ Is the length of the train the same at rest and moving?



Lorentz contraction





$$t_{1} = (L + vt_{1})/c \rightarrow t_{1} = \frac{L}{c - v}$$

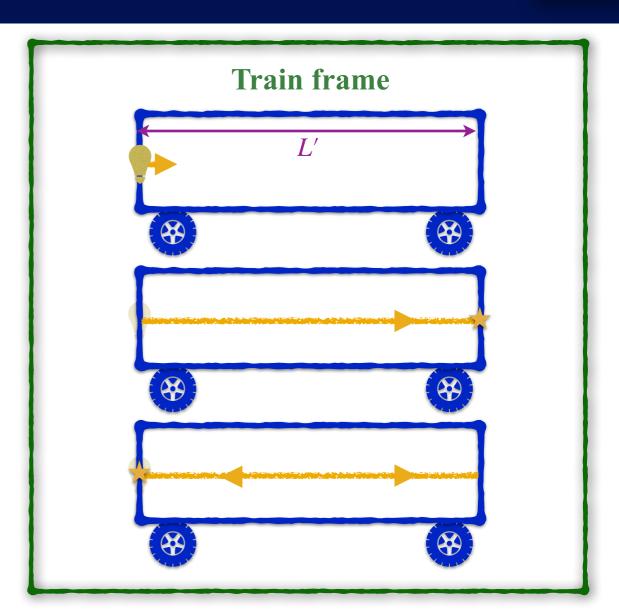
$$t_{2} = (L - vt_{2})/c \rightarrow t_{2} = \frac{L}{c + v}$$

$$t_{1} = \frac{L}{c - v}$$

$$t_{2} = (L - vt_{2})/c \rightarrow t_{2} = \frac{L}{c + v}$$

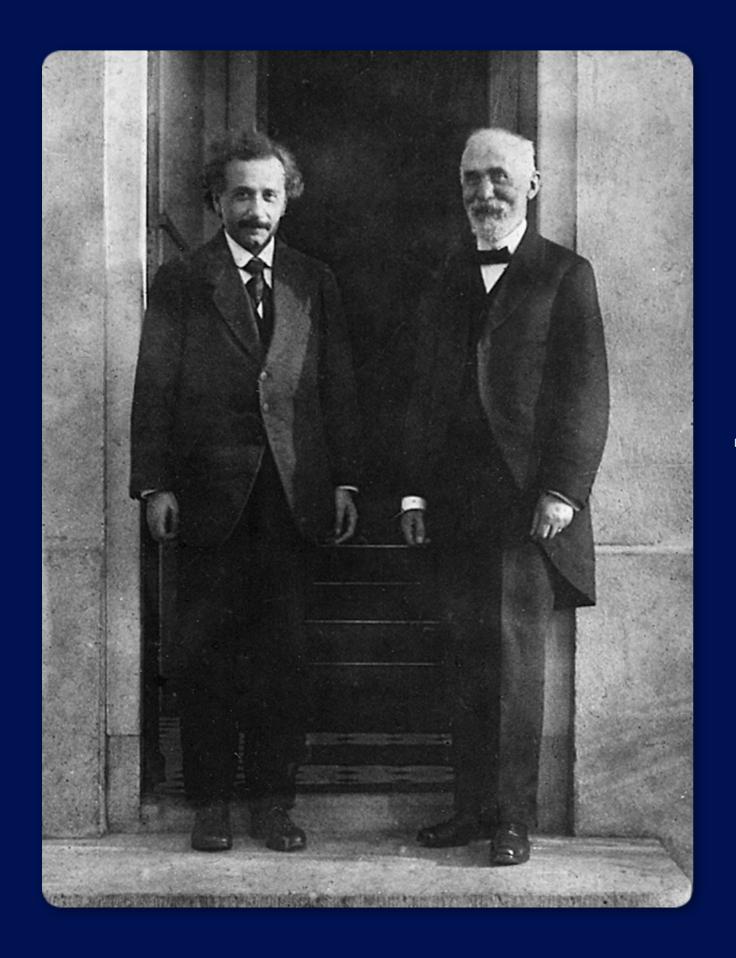
$$t_{3} = t_{1} + t_{2} = \frac{2L}{c} \frac{1}{1 - v^{2}/c^{2}}$$

Using time dilation $t = t'\gamma$



$$t' = \frac{2L'}{c}$$

Moving objects are contracted along the direction of movement



Lorentz transformations and 4-vectors



Lorentz transformations



~ Galilean transformations do not account for time

dilation or space contraction



$$t'=t$$

$$t'=t$$
 $x'=x-vt$ $y'=y$ $z'=z$

$$y' = y$$

$$z'=z$$

Lorentz transformations do

Space and time are $t' = \gamma t - \beta \gamma x/c$ y' = y timately intertwined $x' = \gamma x - \beta \gamma ct$ z' = zintimately intertwined

$$t' = \gamma t - \beta \gamma x / c$$

$$y' = y$$

$$z' = z$$



→ Where we defined

Normalized velocity
$$\beta = \frac{v}{c}$$
,

unitless number between 0 and 1

Lorentz factor
$$\gamma = \frac{1}{\sqrt{1-\beta^2}}$$
,

unitless number between 1 and ∞



4-vectors

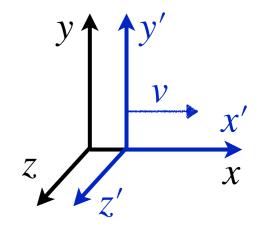


Instead of describing positions with 3D vectors, we use 4-vectors that combine time and space



$$\begin{bmatrix} ct' \\ x' \\ y' \\ z' \end{bmatrix} = \begin{bmatrix} \gamma & -\beta\gamma & 0 & 0 \\ -\beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} ct \\ x \\ y \\ z \end{bmatrix}$$

$$\begin{bmatrix} ct \\ x \\ y \\ z \end{bmatrix}$$
Lorentz transform



Lorentz transformation in matrix form

Normalized velocity $\beta = \frac{v}{c}$, unitless number between 0 and 1

Lorentz factor
$$\gamma = \frac{1}{\sqrt{1-\beta^2}}$$
,

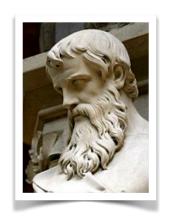
unitless number between 1 and ∞



Norm of 4-vectors



In Euclidean space, norm of vector found adding up the squares of all the components



$$|\vec{r}|^2 = x^2 + y^2 + z^2$$

- ~ 4-vectors live in Minkowski space, and the norm adds the time and spatial components with different signs
 - → This **norm** is the **same in all frames** (invariant)

$$\sum r_{\mu}r^{\mu} = c^2t^2 - x^2 - y^2 - z^2 = c^2\tau^2$$



The proper time is the time that passed for a traveler that followed the space-time trajectory given by r^{μ}

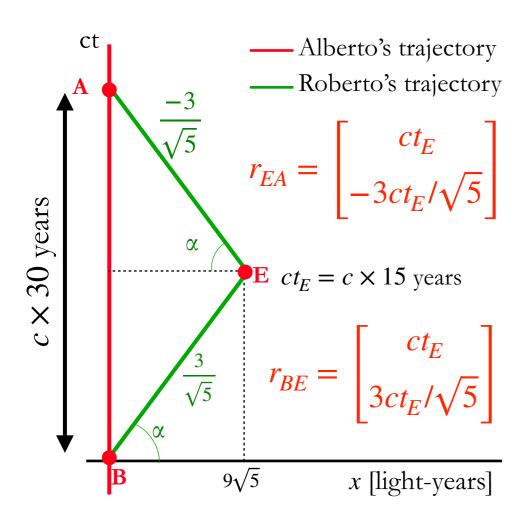


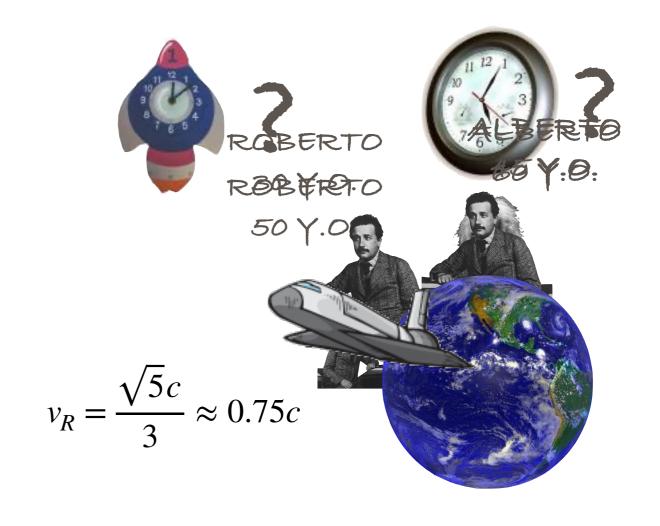
Twin paradox: what's Roberto's age?





$$v_R = \frac{-\sqrt{5}c}{3} \approx -0.75c$$





From B to E, the norm of the vector is

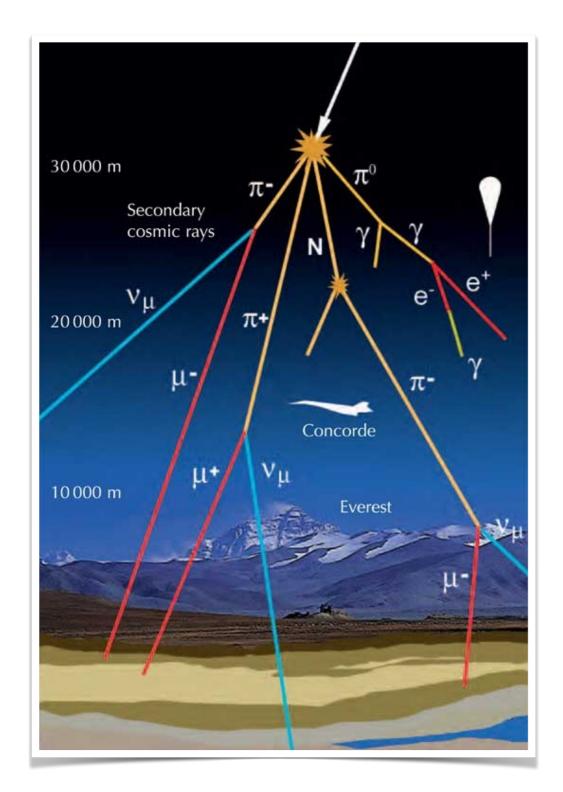
$$|r_{EA}|^2 = |r_{BE}|^2 = c^2 (10 \text{ years})^2$$

Roberto is 50 when he comes back!



Example: lifetime of the muon





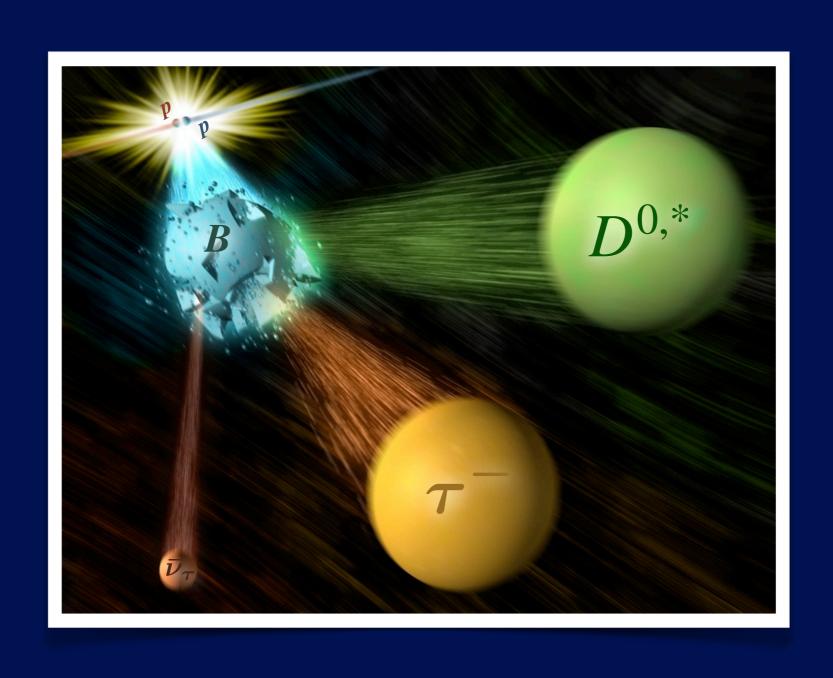
- ~ Atmospheric muons typically come from $\pi^{\pm} \to \mu^{\pm} \nu_{\mu}$ decays
 - → Then decay as $\mu^{\pm} \rightarrow e^{\pm} \nu_e \nu_{\mu}$
- ➤ What is the lifetime of a muon
 - → produced at 25,000 m altitude
 - → traveled vertically towards the Earth
 - \rightarrow decayed at 5,000 m after 66.7 μ s

$$c^2\tau^2 = c^2t^2 - x^2 - y^2 - z^2$$

$$\tau = \frac{\sqrt{c^2 t^2 - x^2}}{c}$$

$$= \frac{\sqrt{c^2(66.7 \times 10^{-6})^2 - 20,000^2}}{c} = 2.2 \ \mu \text{S}$$

Relativistic kinematics: conservation of 4-momentum





The 4-velocity



~ 4-velocity defined as
$$\eta^{\mu} \equiv \frac{dx^{\mu}}{d\tau} = \gamma \frac{dx^{\mu}}{dt} = \gamma(c, v_x, v_y, v_z)$$

 \rightarrow Greek indices like μ refer to the 0, 1, 2, 3 components

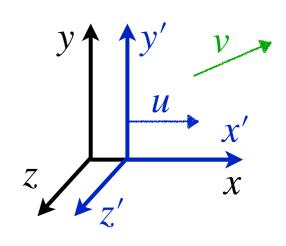
$$\star x^0 = ct, x^1 = x, x^2 = y, x^3 = z$$

Like all 4-vectors

→ Lorentz transforms into other reference frames

$$\eta'^{\mu} = \begin{bmatrix} \gamma(v')c \\ \gamma(v')v'_x \\ \gamma(v')v'_y \\ \gamma(v')v'_z \end{bmatrix} = \begin{bmatrix} \gamma(u) & -\beta(u)\gamma(u) & 0 & 0 \\ -\beta(u)\gamma(u) & \gamma(u) & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \gamma(v)c \\ \gamma(v)v_x \\ \gamma(v)v_y \\ \gamma(v)v_z \end{bmatrix}$$

$$\gamma(v) = \frac{1}{\sqrt{2}}$$



$$\gamma(v) = \frac{1}{\sqrt{1 - v^2/c^2}}$$

→ Its norm is frame invariant

$$\sum \eta_{\mu} \eta^{\mu} = (\gamma c)^{2} - (\gamma v_{x})^{2} - (\gamma v_{y})^{2} - (\gamma v_{y})^{2} = \gamma^{2} (c^{2} - v^{2}) = c^{2}$$



4-momentum

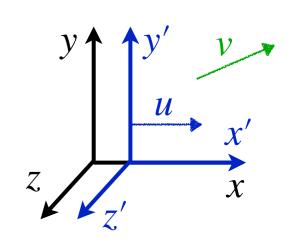


- ~ 4-momentum defined as $p^{\mu} = m\eta^{\mu} = \left(m\gamma c, m\gamma \overrightarrow{v}\right)$
 - \rightarrow This is for a particle of mass m
 - → Also defined for a **set of particles** as the **sum of individual 4-momenta**

∼ Like all 4-vectors

→ Lorentz transforms into other reference frames

$$p'^{\mu} = \begin{bmatrix} m\gamma(v')c \\ m\gamma(v')v'_x \\ m\gamma(v')v'_y \\ m\gamma(v')v'_z \end{bmatrix} = \begin{bmatrix} E'/c \\ p'_x \\ p'_y \\ p'_z \end{bmatrix} = \begin{bmatrix} \gamma(u) & -\beta(u)\gamma(u) & 0 & 0 \\ -\beta(u)\gamma(u) & \gamma(u) & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} E/c \\ p_x \\ p_y \\ p_z \end{bmatrix}$$



→ Its norm is frame invariant

$$\sum p_{\mu}p^{\mu} = E^2/c^2 - p_x^2 - p_y^2 - p_z^2 = m^2c^2 \rightarrow E = \sqrt{m^2c^4 - p^2c^2}$$

$$p^2 = p_x^2 + p_y^2 + p_z^2$$



4-momentum for $v \ll c$



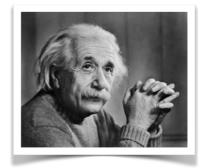
~ Spatial part looks like standard momentum

$$p^{\mu} = \begin{bmatrix} m\gamma c \\ m\gamma v_x \\ m\gamma v_y \\ m\gamma v_z \end{bmatrix} = \begin{bmatrix} E/c \\ p_x \\ p_y \\ p_z \end{bmatrix} \quad m\gamma \overrightarrow{v} = \frac{m\overrightarrow{v}}{\sqrt{1 - \frac{v^2}{c^2}}} = \overrightarrow{m} \overrightarrow{v} + m\overrightarrow{v} \frac{v^2}{2c^2} + \mathcal{O}\left(\frac{v^4}{c^4}\right)$$

 \sim Time part is mc^2 constant plus kinetic energy

$$\frac{E}{c} = m\gamma c = \frac{mc}{\sqrt{1 - \frac{v^2}{c^2}}} = \frac{1}{c} \left[mc^2 + m\frac{v^2}{2} + \mathcal{O}\left(\frac{v^4}{c^2}\right) \right]$$

Einsteins brilliant insight $E=mc^2$ (at rest)



When not at rest,

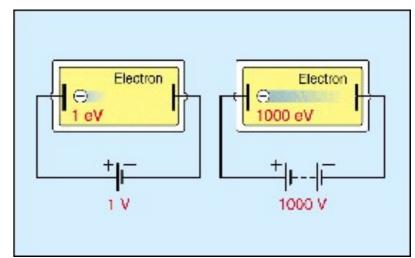
$$E = \sqrt{m^2c^4 - p^2c^2}$$



Units in particle physics



- ~ Typically use *electronvolts* (eV) to measure energy
 - → Energy of an electron accelerated by a 1 V potential
 - → Very useful in the original experiments
 - \rightarrow 1 eV = 1.602×10⁻¹⁹ J
- ~ Use eV/c for momentum



∼ What is the energy of an electron ($m_e = 511 \text{ keV/c}^2$) that has a momentum of 1 MeV/c?

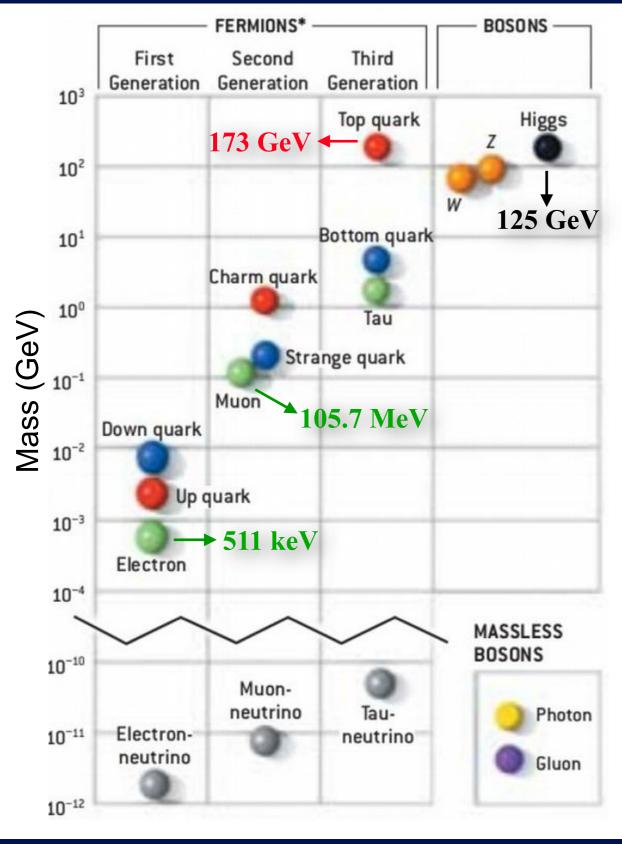
$$E_e = \sqrt{p_e^2 c^2 + m_e^2 c^4} = \sqrt{(1 \text{ MeV/}c)^2 c^2 + (0.511 \text{ MeV/}c^2)^2 c^4} = 1.12 \text{ MeV}$$

Very often, we just set c = 1, so energy, mass, and momentum are measured in eV



Particle masses





- Elementary particles masses cover a large range
 - → Photon/gluon are massless
 - Neutrinos ~ meV
 - → 1st generation fermions ~1 MeV
 - → 2nd generation fermions ~ 0.1 1 GeV
 - → 3rd generation fermions ~ 1 173 GeV
 - → W, Z bosons ~ 90 GeV
 - → Higgs boson ~ 125 GeV
- Masses of non-elementary particles heavily determined by binding energy
 - \Rightarrow eg, m(u) = 2.3 MeV, m(d) = 4.8 MeV
 - + Mass of proton (uud) is 938 MeV



4-momentum conservation



- → In SR, it is 4-momentum that is conserved
 - → Generalization of energy and 3-momentum conservation
- \sim For a closed system of n particles

$$\sum_{i}^{n} p_{\mu}^{i,ini} = \sum_{i}^{n} p_{\mu}^{i,fin} \quad \text{for each coordinate } \mu$$

$$\begin{bmatrix} \frac{E^{1,ini}}{c} \\ p_{x}^{1,ini} \\ p_{y}^{1,ini} \\ p_{z}^{1,ini} \end{bmatrix} + \begin{bmatrix} \frac{E^{2,ini}}{c} \\ p_{x}^{2,ini} \\ p_{z}^{2,ini} \\ p_{z}^{2,ini} \end{bmatrix} + \dots + \begin{bmatrix} \frac{E^{n,ini}}{c} \\ p_{x}^{n,ini} \\ p_{y}^{n,ini} \\ p_{z}^{1,fin} \end{bmatrix} = \begin{bmatrix} \frac{E^{1,fin}}{c} \\ p_{x}^{1,fin} \\ p_{y}^{1,fin} \\ p_{z}^{1,fin} \end{bmatrix} + \begin{bmatrix} \frac{E^{2,fin}}{c} \\ p_{x}^{2,fin} \\ p_{y}^{2,fin} \\ p_{z}^{2,fin} \end{bmatrix} + \dots + \begin{bmatrix} \frac{E^{n,fin}}{c} \\ p_{x}^{n,fin} \\ p_{y}^{n,fin} \\ p_{z}^{n,fin} \end{bmatrix}$$



Degrees of freedom in particle decay



 \sim A particle α decays into two other particles λ and ω : $\alpha \to \lambda \omega$

General conservation of 4-momentum

$$\begin{bmatrix} \sqrt{p_{\alpha}^{2} + m_{\alpha}^{2}c^{2}} \\ p_{\alpha,x} \\ p_{\alpha,y} \\ p_{\alpha,z} \end{bmatrix} = \begin{bmatrix} \sqrt{p_{\lambda}^{2} + m_{\lambda}^{2}c^{2}} \\ p_{\lambda,x} \\ p_{\lambda,y} \\ p_{\lambda,z} \end{bmatrix} + \begin{bmatrix} \sqrt{p_{\omega}^{2} + m_{\omega}^{2}c^{2}} \\ p_{\omega,x} \\ p_{\omega,y} \\ p_{\omega,z} \end{bmatrix} = \begin{bmatrix} \sqrt{p_{\lambda}^{2} + m_{\lambda}^{2}c^{2}} \\ p_{\lambda} \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} \sqrt{p_{\lambda}^{2} + m_{\lambda}^{2}c^{2}} \\ p_{\lambda} \\ 0 \\ 0 \end{bmatrix} + \begin{bmatrix} \sqrt{p_{\omega}^{2} + m_{\omega}^{2}c^{2}} \\ p_{\omega} \\ 0 \\ 0 \end{bmatrix}$$

Conservation of 4-momentum in two-body decay

$$\begin{bmatrix} m_{\alpha}c \\ 0 \\ 0 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} \sqrt{p_{\lambda}^2 + m_{\lambda}^2 c^2} \\ p_{\lambda} \\ 0 \\ 0 \end{bmatrix} + \begin{bmatrix} \sqrt{p_{\omega}^2 + m_{\omega}^2 c^2} \\ p_{\omega} \\ 0 \\ 0 \end{bmatrix}$$

$$p_{\alpha} = \sqrt{p_{\alpha,x}^2 + p_{\alpha,y}^2 + p_{\alpha,x}^z}$$

$$p_{\lambda} = \sqrt{p_{\lambda,x}^2 + p_{\lambda,y}^2 + p_{\lambda,x}^z}$$

$$p_{\omega} = \sqrt{p_{\omega,x}^2 + p_{\omega,y}^2 + p_{\omega,x}^z}$$

	Degrees of freedom				Equations
	α	λ	ω	Total	(constraints)
Original	4	4	4	12	4
Masses	3	3	3	9	4
α at rest	0	3	3	6	4
Symmetry	0	1	1	2	2

$$\begin{array}{cccc} p_{\lambda}^{\nu} & p_{\alpha}^{\nu} & p_{\omega}^{\nu} \\ & & \end{array}$$

Cannot know λ and ω direction (isotropic decay), but we know that they will **both** travel along the same line in opposite directions by conservation of momentum → Choose x axis along the direction of the decay



(Forbidden) decay of a muon



~ What is the energy of the (massless) photon γ emitted in the decay $\mu^- \to e^- + \gamma$ assuming the muon is initially

at rest? p_{γ}^{ν} p_{μ}^{ν} p_{e}^{ν}

Cannot know γ and e-direction (isotropic decay), but we know that they will both travel along the same line in opposite directions by conservation of momentum

Since the muon is at rest, $p_{\mu} = 0$, the photon is massless, $m_{\gamma} = 0$

$$\begin{bmatrix} m_{\mu}c \\ 0 \end{bmatrix} = \begin{bmatrix} \sqrt{p_e^2 + m_e^2c^2} \\ p_e \end{bmatrix} + \begin{bmatrix} p_{\gamma} \\ p_{\gamma} \end{bmatrix} \Rightarrow \sqrt{p_e^2 + m_e^2c^2} = m_{\mu}c - p_{\gamma}$$

$$0 = p_e + p_{\gamma}$$

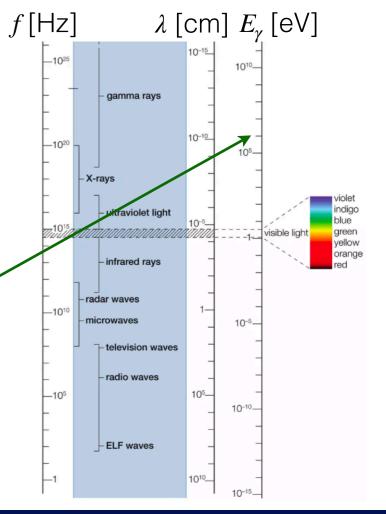
Solve for positive p_{γ}

$$p_{\gamma}^{2} + m_{e}^{2}c^{2} = m_{\mu}^{2}c^{2} - 2m_{\mu}cp_{\gamma} + p_{\gamma}^{2} \Rightarrow p_{\gamma} = \frac{m_{\mu}^{2}c - m_{e}^{2}c}{2m_{\mu}}$$

$$E_{\gamma} = p_{\gamma}c = \frac{m_{\mu}^{2}c^{2} - m_{e}^{2}c^{2}}{2m_{\mu}} = 52.83 \text{ MeV}$$

Conservation of 4-momentum

$$\begin{bmatrix} \sqrt{p_{\mu}^2 + m_{\mu}^2 c^2} \\ p_{\mu} \end{bmatrix} = \begin{bmatrix} \sqrt{p_e^2 + m_e^2 c^2} \\ p_e \end{bmatrix} + \begin{bmatrix} \sqrt{p_{\gamma}^2 + m_{\gamma}^2 c^2} \\ p_{\gamma} \end{bmatrix}$$

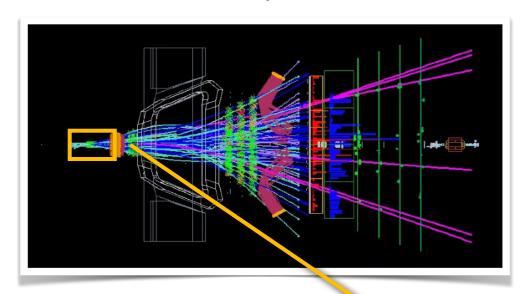




Reconstructing a particle

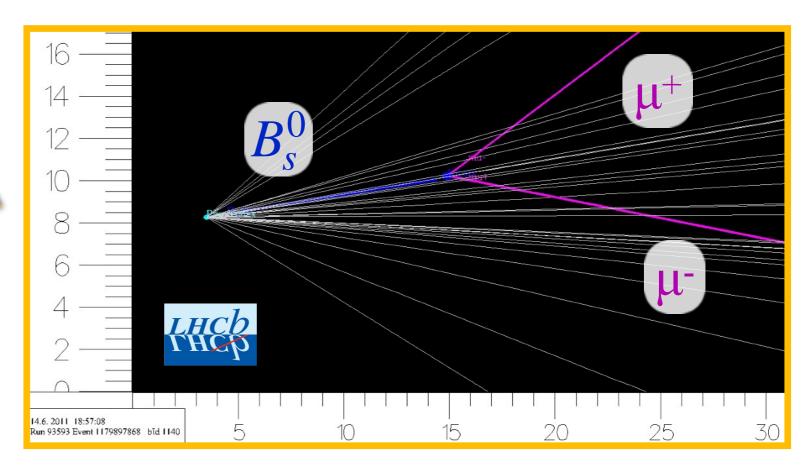


- ∼ How do you discover a new particle?
 - → Look for particles decaying from a common vertex



$$B_s^0 \to \mu^+\mu^-$$

The B_s^0 is invisible, but we see the muons and reconstruct its momentum





Reconstruct invariant mass



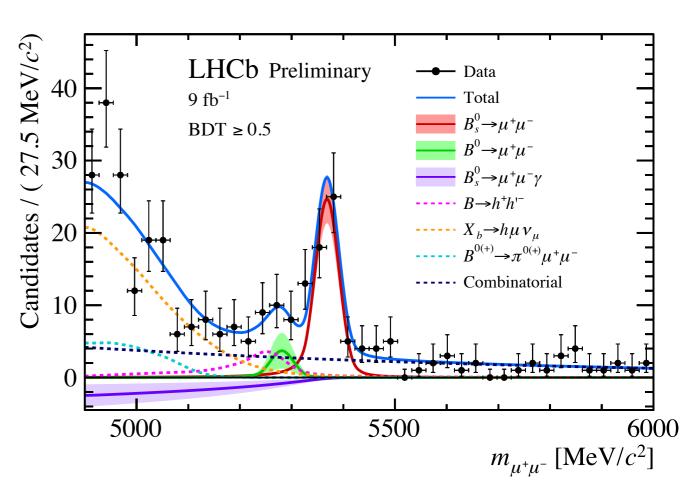
- How do you discover a new particle?
 - → Look for particles decaying from a common vertex
 - → Add up the momenta of the daughter particles

$$\begin{bmatrix} \sqrt{p_1^2 + m_{\mu}^2 c^2} \\ p_1 \end{bmatrix} + \begin{bmatrix} \sqrt{p_2^2 + m_{\mu}^2 c^2} \\ p_2 \end{bmatrix} = \begin{bmatrix} E_B \\ p_B \end{bmatrix}$$

→ Calculate invariant mass

$$m_B = \sqrt{E_B^2 - p_B^2}$$

→ Plot histogram and find peak!





Particle accelerator types



Geometry

Linear collider
Stanford Linear
Collider (SLC)



VS



Circular collider

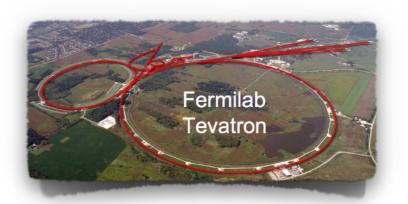
Large ElectronPositron collider (LEP)

Type of collision

Fixed target NA62



VS



Collider *Tevatron*

Type of particle

Lepton collider
Positron-Electron
Project II (PEP-II)



VS



Hadron collider

Large Hadron

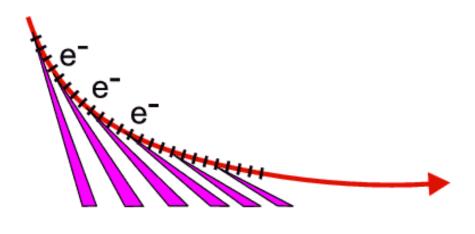
Collider (LHC)



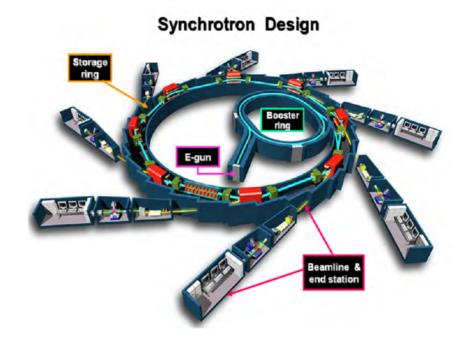
Linear vs Circular colliders



- ~ Circular geometry has key advantages as particles circle around
 - → Can have higher energies as particles are accelerated at each cycle
 - → Can have **higher collision rates** as particle bunches have more opportunities to collide
- However, charged particles emit synchrotron radiation when changing direction
 - → Power loss grows very quickly with energy
 - → Circular geometry **not feasible for light particles** (eg, electrons) for $\sqrt{s} \ge \sim 250$ GeV
 - ◆ Use protons (m_p ~ 2000m_e) for higher energies



$$P_{\rm sync}^{\rm loss} \propto \frac{E^4}{m^4 R^2}$$





Fixed target vs Colliders



In Fixed target experiments a beam of high-energy particles collides against a large and stationary target

$$p_1^{\nu}$$
 p_3^{ν} p_2^{ν}

Ep1
$$p_{p1} = \sqrt{E_{p1}^2 - m_p^2}$$
 $p_{p2} = (1 \text{ GeV/c}, 0, 0, 0)$ $p_{p1} = (10 \text{ TeV/c}, 10 \text{ TeV/c}, 0, 0)$ $p_{p1} = (10 \text{ TeV/c})$ Conservation of 4-momentum

 $\left| \sqrt{p_1^2 + m_1^2 c^2} \right| + \left| \sqrt{p_2^2 + m_2^2 c^2} \right| = \left| \sqrt{p_3^2 + m_3^2 c^2} \right|$

$$\begin{bmatrix} p_1 \\ p_1 \end{bmatrix} + \begin{bmatrix} m_p \\ 0 \end{bmatrix} = \begin{bmatrix} \sqrt{p_3^2 + m_3^2 c^2} \\ p_3 \end{bmatrix} \qquad p_3 = p_1$$

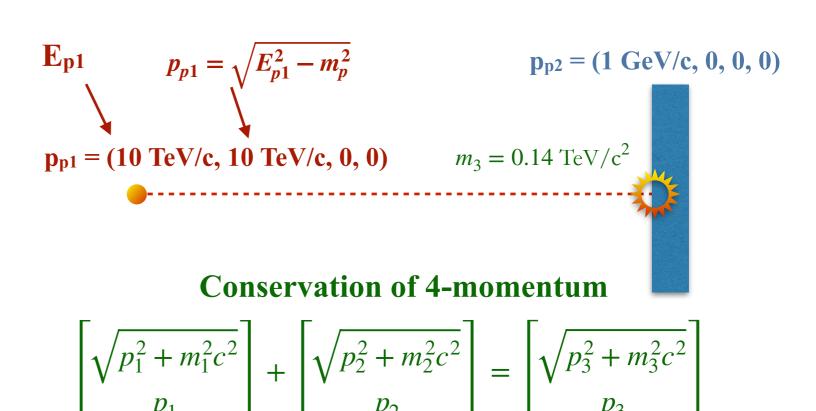
$$m_3^2 c^2 = \left(p_1 + m_p c \right)^2 - p_1^2 = 2p_1 m_p c \quad \rightarrow \quad m_3 = 0.14 \text{ TeV/c}^2$$



Fixed target vs Colliders



- In Fixed target experiments a beam of high-energy particles collides against a large and stationary target
 - → Easy, cheap
 - Inefficient → difficult to reach high energies
- Colliders crash two narrow beams of particles
 - → Difficult but efficient



$$p_{p1} = (10 \text{ TeV/c}, 10 \text{ TeV/c}, 0, 0)$$
 $p_{p2} = (10 \text{ TeV/c}, -10 \text{ TeV/c}, 0, 0)$
 $m_3 = 20 \text{ TeV/c}^2$



Summary of 4-momentum conservation



Conservation of 4-momentum allows us to solve particle physics problems

Conservation of 4-momentum in a particle decay

2-body decays are always collinear

Conservation of 4-momentum in a head on particle collision

$$p_1^{\nu}$$
 p_3^{ν} p_2^{ν}

$$\begin{bmatrix}
p_1^{\nu} & p_3^{\nu} & p_2^{\nu} \\
p_1 & p_2
\end{bmatrix} + \begin{bmatrix}
\sqrt{p_1^2 + m_1^2 c^2} \\
p_2 & p_2
\end{bmatrix} + \begin{bmatrix}
\sqrt{p_2^2 + m_2^2 c^2} \\
p_2 & p_3
\end{bmatrix} = \begin{bmatrix}
\sqrt{p_3^2 + m_3^2 c^2} \\
p_3 & p_3
\end{bmatrix}$$

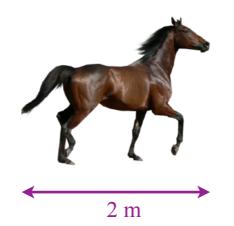
Backup

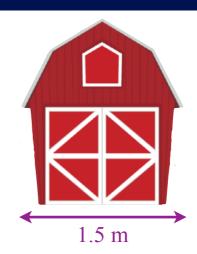


Horse-barn paradox



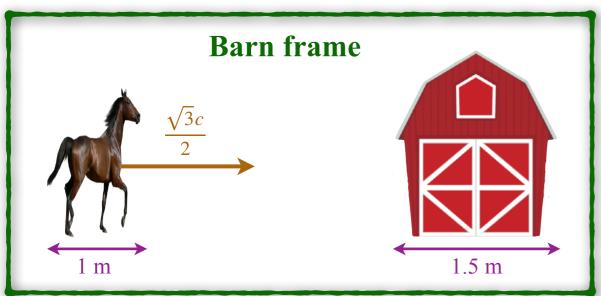
Could we **fit** a **2 m horse** into a **1.5 m barn** using relativity?



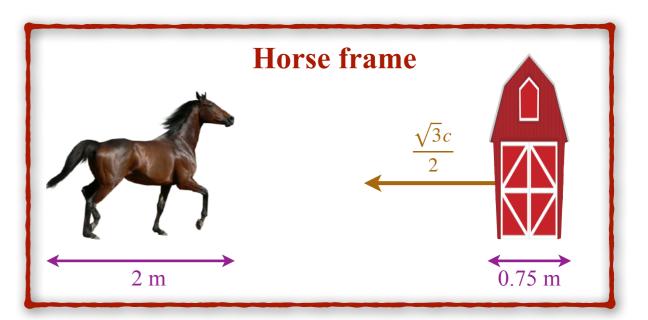


Yes! Accelerate the horse to

$$v_{\text{horse}} = \frac{\sqrt{3}c}{2} \rightarrow \gamma = \frac{1}{\sqrt{1 - \frac{v_{\text{horse}}^2}{c^2}}} = 2$$



How can the 2 m horse both fit in the 1.5 m barn in the rest frame and not fit in the barn in its own frame?



Lack of simultaneity is what leads to many of SR's paradoxes

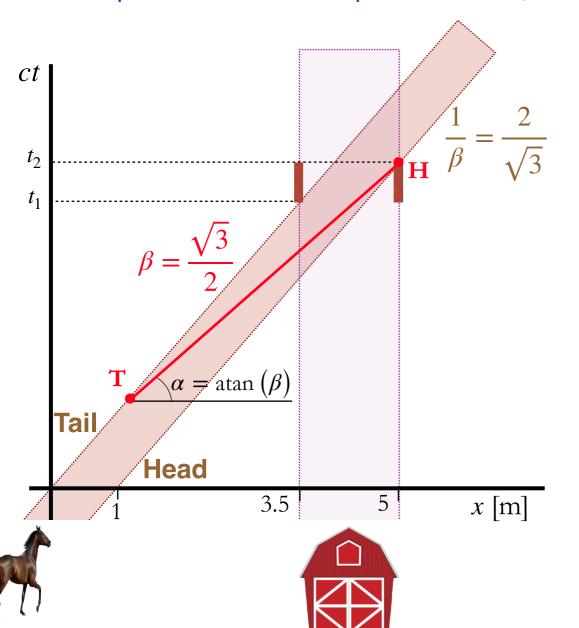


Horse-barn paradox resolution



→ It is simply a problem of simultaneity

→ The tail and head of the horse being inside the barn at the same time depends on the speed that you are going



In the barn rest frame, the horse is $(2m)/\gamma = 1m$ long, and the slopes of the head and tail of the horse are $\frac{1}{\beta} \approx 1.15$ Since the barn is at rest in this frame, its doors have infinite slope.

Between t₁ and t₂ the horse is fully enclosed inside the barn because both its tail and head are inside.



Quantum mechanics in a nutshell



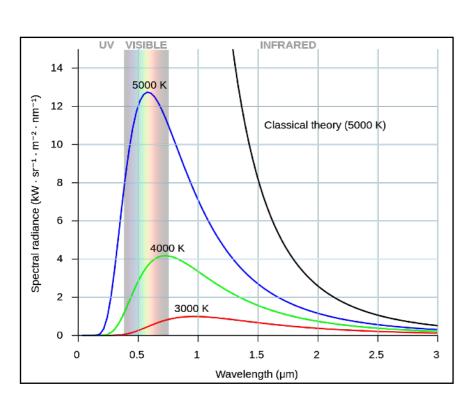
Black body radiation Ultraviolet catastrophe

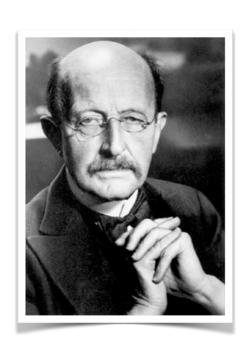


Energy quantization

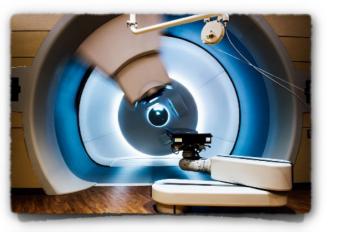


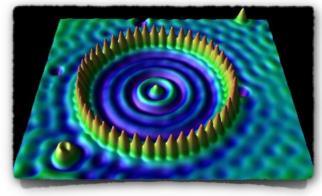
Deeper understanding of the Universe's fundamental constituents













+ many other experiments (photoelectric effect...)

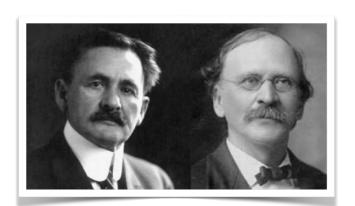
+ work from many others (Einstein, Schrödinger, Heisenberg...)

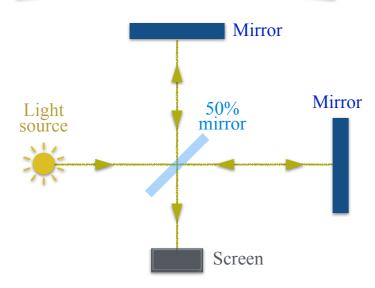


Relativity in a nutshell



Michelson-Morley
experiment
Could not find ether, the
medium light travels
with respect to

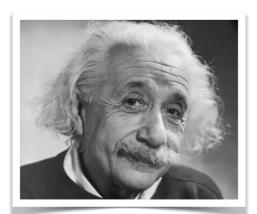




+ experiments from many others (ether dragging...)



Einstein's postulates



The laws of Physics have the same form in all inertial frames of reference

The speed of light in vacuum has the same value c in every direction in all inertial frames of reference

> + work from many others (Poincaré, Lorentz...)

Deeper understanding of the Universe's spacetime, mass, and energy

$$E = mc^2$$



$$\Delta t' = \frac{\Delta t}{\sqrt{1 - \frac{v^2}{c^2}}}$$

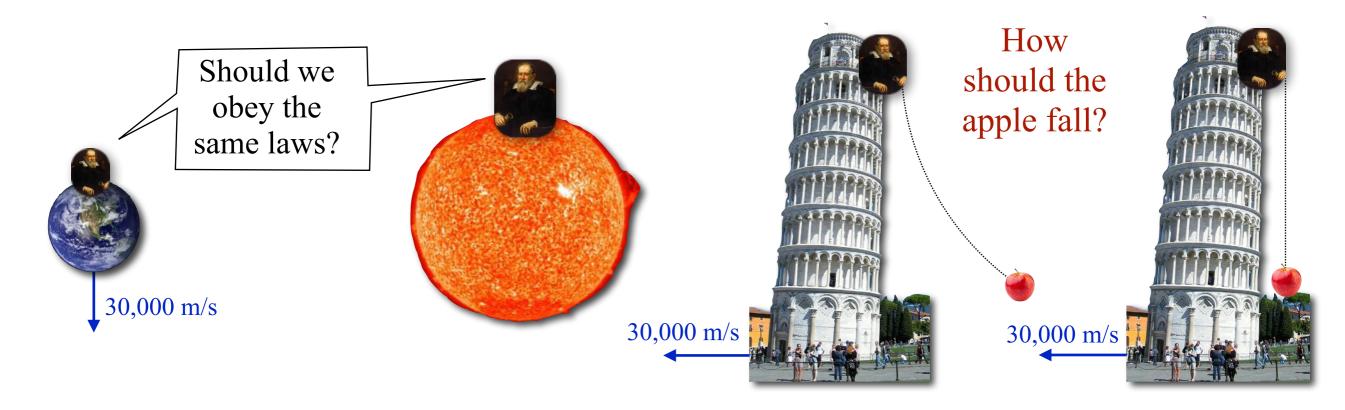




Galilean relativity



 Not obvious whether Earth's laws of physics should apply elsewhere



Postulate from Galileo Galileo (1564 - 1642)

The laws of Motion have the same form in all inertial frames of reference

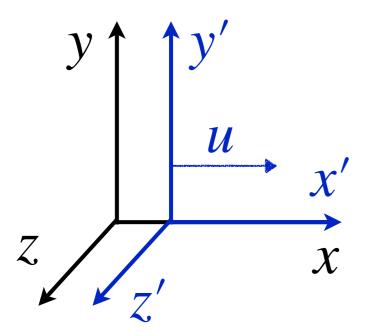
Inertial frames of reference are those that move at constant relative velocities



Galilean transformations



- ∼ The Galilean transformations rely on two axioms
 - → There exists an absolute space
 - → All inertial frames share a universal time
- Thus, the coordinate transformation between two inertial frames where the origins are the same at t=0 and the velocity u is aligned with the x axis



$$t' = t \qquad x' = x - ut$$
$$y' = y \qquad z' = z$$



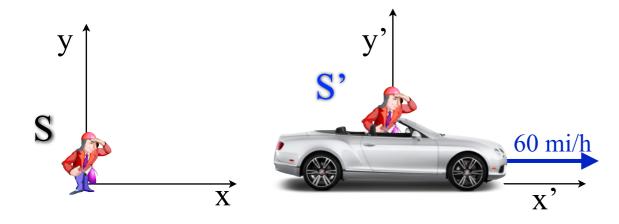


Galilean example



 \sim For two frames of references S and S' with the same origin at t = 0 and aligned

$$t' = t x' = x - (60 \text{ mi/h}) t$$
$$y' = y z' = z$$





$$x_{\text{train}} = 10 + (200 \,\text{mi/h})t - (10 \,\text{mi/h}^2)t^2$$

$$L_{\text{train}} = 20 \,\text{m}$$

$$v_{\rm train} =$$

$$a_{\rm train} =$$

$$x'_{\text{train}} =$$

$$L'_{\text{train}} =$$

$$v'_{\text{train}} =$$

$$a'_{\text{train}} =$$